## Solution to Theoretical Question 1

## A Swing with a Falling Weight

### Part A

(a) Since the length of the string  $L = s + R\theta$  is constant, its rate of change must be zero. Hence we have

$$\dot{s} + R\dot{\theta} = 0 \tag{A1}$$

(b) Relative to O, Q moves on a circle of radius R with angular velocity  $\dot{\theta}$ , so

$$\vec{v}_O = R\dot{\theta}\,\hat{t} = -\dot{s}\,\hat{t} \tag{A2}$$

(c) Refer to Fig. A1. Relative to Q, the displacement of P in a time interval  $\Delta t$  is  $\Delta \vec{r}' = (s\Delta\theta)(-\hat{r}) + (\Delta s)\hat{t} = [(s\dot{\theta})(-\hat{r}) + \dot{s}\hat{t}]\Delta t$ . It follows

$$\vec{v}' = -s \dot{\theta} \, \hat{r} + \dot{s} \, \hat{t} \tag{A3}$$

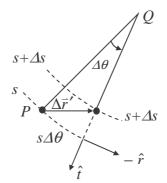


Figure A1

(d) The velocity of the particle relative to *O* is the sum of the two relative velocities given in Eqs. (A2) and (A3) so that

$$\vec{v} = \vec{v}' + \vec{v}_Q = (-s\dot{\theta}\,\hat{r} + \dot{s}\,\hat{t}\,) + R\dot{\theta}\,\hat{t} = -s\dot{\theta}\,\hat{r} \tag{A4}$$

(e) Refer to Fig. A2. The  $(-\hat{t})$ -component of the velocity change  $\Delta \vec{v}$  is given by  $(-\hat{t}) \cdot \Delta \vec{v} = v \Delta \theta = v \dot{\theta} \Delta t$ . Therefore, the  $\hat{t}$ -component of the acceleration  $\vec{a} = \Delta \vec{v} / \Delta t$  is given by  $\hat{t} \cdot \hat{a} = -v \dot{\theta}$ . Since the speed v of the particle is  $s \dot{\theta}$  according to Eq. (A4), we see that the  $\hat{t}$ -component of the particle's acceleration at P is given by



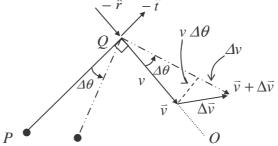
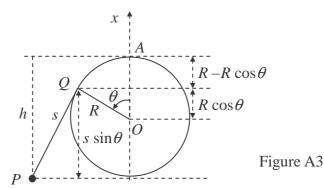


Figure A2

Note that, from Fig. A2, the radial component of the acceleration may also be obtained as  $\vec{a} \cdot \hat{r} = -dv/dt = -d(s\dot{\theta})/dt$ .

(f) Refer to Fig. A3. The gravitational potential energy of the particle is given by U = -mgh. It may be expressed in terms of s and  $\theta$  as

$$U(\theta) = -mg[R(1-\cos\theta) + s\sin\theta] \tag{A6}$$



(g) At the lowest point of its trajectory, the particle's gravitational potential energy U must assume its minimum value  $U_m$ . By differentiating Eq. (A6) with respect to  $\theta$  and using Eq. (A1), the angle  $\theta_m$  corresponding to the minimum gravitational energy can be obtained.

$$\frac{dU}{d\theta} = -mg\left(R\sin\theta + \frac{ds}{d\theta}\sin\theta + s\cos\theta\right)$$
$$= -mg\left[R\sin\theta + (-R)\sin\theta + s\cos\theta\right]$$
$$= -mgs\cos\theta$$

At  $\theta = \theta_m$ ,  $\frac{dU}{d\theta}\Big|_{\theta_m} = 0$ . We have  $\theta_m = \frac{\pi}{2}$ . The lowest point of the particle's trajectory is

shown in Fig. A4 where the length of the string segment of QP is  $s = L - \pi R/2$ .

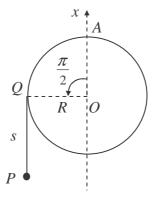


Figure A4

From Fig. A4 or Eq. (A6), the minimum potential energy is then

$$U_{m} = U(\pi/2) = -mg[R + L - (\pi R/2)]$$
(A7)

Initially, the total mechanical energy E is 0. Since E is conserved, the speed  $v_m$  of the particle at the lowest point of its trajectory must satisfy

$$E = 0 = \frac{1}{2}mv_m^2 + U_m \tag{A8}$$

From Eqs. (A7) and (A8), we obtain

$$v_m = \sqrt{-2U_m/m} = \sqrt{2g[R + (L - \pi R/2)]}$$
 (A9)

## Part B

(h) From Eq. (A6), the total mechanical energy of the particle may be written as

$$E = 0 = \frac{1}{2}mv^{2} + U(\theta) = \frac{1}{2}mv^{2} - mg[R(1 - \cos\theta) + s\sin\theta]$$
 (B1)

From Eq. (A4), the speed v is equal to  $s\dot{\theta}$ . Therefore, Eq. (B1) implies

$$v^{2} = (s\dot{\theta})^{2} = 2g[R(1-\cos\theta) + s\sin\theta]$$
 (B2)

Let T be the tension in the string. Then, as Fig. B1 shows, the  $\hat{t}$ -component of the net force on the particle is  $-T + mg \sin \theta$ . From Eq. (A5), the tangential acceleration of the particle is  $(-s\dot{\theta}^2)$ . Thus, by Newton's second law, we have

$$m(-s\dot{\theta}^2) = -T + mg\sin\theta \tag{B3}$$

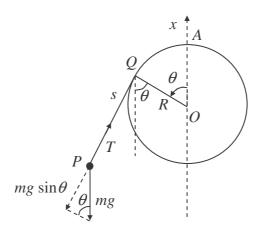


Figure B1

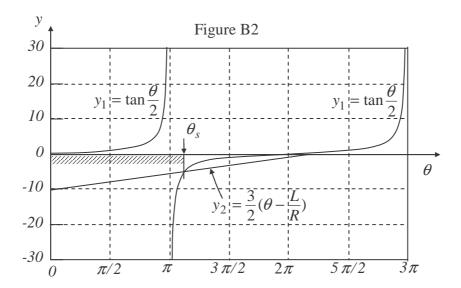
According to the last two equations, the tension may be expressed as

$$T = m(s\dot{\theta}^{2} + g\sin\theta) = \frac{mg}{s} [2R(1-\cos\theta) + 3s\sin\theta]$$

$$= \frac{2mgR}{s} [\tan\frac{\theta}{2} - \frac{3}{2}(\theta - \frac{L}{R})](\sin\theta)$$

$$= \frac{2mgR}{s} (y_{1} - y_{2})(\sin\theta)$$
(B4)

The functions  $y_1 = \tan(\theta/2)$  and  $y_2 = 3(\theta - L/R)/2$  are plotted in Fig B2.



From Eq. (B4) and Fig. B2, we obtain the result shown in Table B1. The angle at which  $y_2 = y_1$  is called  $\theta_s$  ( $\pi < \theta_s < 2\pi$ ) and is given by

$$\frac{3}{2}(\theta_s - \frac{L}{R}) = \tan\frac{\theta_s}{2} \tag{B5}$$

or, equivalently, by

$$\frac{L}{R} = \theta_s - \frac{2}{3} \tan \frac{\theta_s}{2} \tag{B6}$$

Since the ratio L/R is known to be given by

$$\frac{L}{R} = \frac{9\pi}{8} + \frac{2}{3}\cot\frac{\pi}{16} = (\pi + \frac{\pi}{8}) - \frac{2}{3}\tan\frac{1}{2}(\pi + \frac{\pi}{8})$$
(B7)

one can readily see from the last two equations that  $\theta_s = 9\pi/8$ .

Table B1

	$(y_1 - y_2)$	$\sin \theta$	tension T
$0 < \theta < \pi$	positive	positive	positive
$\theta = \pi$	+ ∞	0	positive
$\pi < \theta < \theta_s$	negative	negative	positive
$\theta = \theta_{s}$	zero	negative	zero
$\theta_s < \theta < 2\pi$	positive	negative	negative

Table B1 shows that the tension T must be positive (or the string must be taut and straight) in the angular range  $0<\theta<\theta_s$ . Once  $\theta$  reaches  $\theta_s$ , the tension T becomes zero and the part of the string not in contact with the rod will not be straight afterwards. The shortest possible value  $s_{\min}$  for the length s of the line segment QP therefore occurs at  $\theta=\theta_s$  and is given by

$$s_{\min} = L - R\theta_s = R(\frac{9\pi}{8} + \frac{2}{3}\cot\frac{\pi}{16} - \frac{9\pi}{8}) = \frac{2R}{3}\cot\frac{\pi}{16} = 3.352R$$
 (B8)

When  $\theta = \theta_s$ , we have T = 0 and Eqs. (B2) and (B3) then leads to  $v_s^2 = -gs_{\min} \sin \theta_s$ . Hence the speed  $v_s$  is

$$v_s = \sqrt{-gs_{\min}\sin\theta_s} = \sqrt{\frac{2gR}{3}\cot\frac{\pi}{16}\sin\frac{\pi}{8}} = \sqrt{\frac{4gR}{3}}\cos\frac{\pi}{16}$$

$$= 1.133\sqrt{gR}$$
(B9)

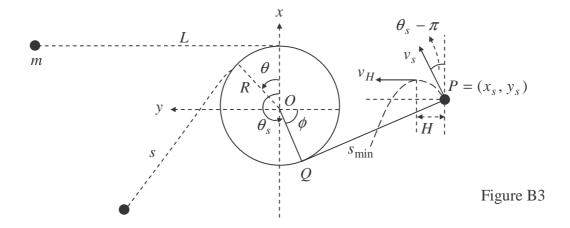
(i) When  $\theta \ge \theta_s$ , the particle moves like a projectile under gravity. As shown in Fig. B3, it is projected with an initial speed  $v_s$  from the position  $P = (x_s, y_s)$  in a direction making an angle  $\phi = (3\pi/2 - \theta_s)$  with the y-axis.

The speed  $v_H$  of the particle at the highest point of its parabolic trajectory is equal to the y-component of its initial velocity when projected. Thus,

$$v_H = v_s \sin(\theta_s - \pi) = \sqrt{\frac{4gR}{3}} \cos\frac{\pi}{16} \sin\frac{\pi}{8} = 0.4334\sqrt{gR}$$
 (B10)

The horizontal distance H traveled by the particle from point P to the point of maximum height is

$$H = \frac{v_s^2 \sin 2(\theta_s - \pi)}{2g} = \frac{v_s^2}{2g} \sin \frac{9\pi}{4} = 0.4535R$$
 (B11)



The coordinates of the particle when  $\theta = \theta_s$  are given by

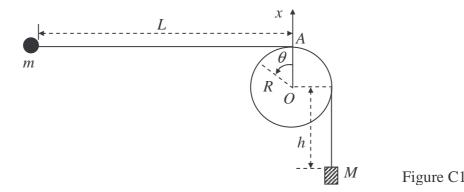
$$x_s = R\cos\theta_s - s_{\min}\sin\theta_s = -R\cos\frac{\pi}{8} + s_{\min}\sin\frac{\pi}{8} = 0.358R$$
 (B12)

$$y_s = R\sin\theta_s + s_{\min}\cos\theta_s = -R\sin\frac{\pi}{8} - s_{\min}\cos\frac{\pi}{8} = -3.478R$$
 (B13)

Evidently, we have  $|y_s| > (R+H)$ . Therefore the particle can indeed reach its maximum height without striking the surface of the rod.

### Part C

(j) Assume the weight is initially lower than O by h as shown in Fig. C1.



When the weight has fallen a distance D and stopped, the law of conservation of total mechanical energy as applied to the particle-weight pair as a system leads to

$$-Mgh = E' - Mg(h+D) \tag{C1}$$

where E' is the *total mechanical energy of the particle* when the weight has stopped. It follows

$$E' = MgD \tag{C2}$$

Let  $\Lambda$  be the total length of the string. Then, its value at  $\theta = 0$  must be the same as at any other angular displacement  $\theta$ . Thus we must have

$$\Lambda = L + \frac{\pi}{2}R + h = s + R(\theta + \frac{\pi}{2}) + (h+D)$$
 (C3)

Noting that  $D = \alpha L$  and introducing  $\ell = L - D$ , we may write

$$\ell = L - D = (1 - \alpha)L \tag{C4}$$

From the last two equations, we obtain

$$s = L - D - R\theta = \ell - R\theta \tag{C5}$$

After the weight has stopped, the total mechanical energy of the particle must be conserved. According to Eq. (C2), we now have, instead of Eq. (B1), the following equation:

$$E' = MgD = \frac{1}{2}mv^2 - mg[R(1 - \cos\theta) + s\sin\theta]$$
 (C6)

The square of the particle's speed is accordingly given by

$$v^{2} = (s\dot{\theta})^{2} = \frac{2MgD}{m} + 2gR \left[ (1 - \cos\theta) + \frac{s}{R}\sin\theta \right]$$
 (C7)

Since Eq. (B3) stills applies, the tension T of the string is given by

$$-T + mg\sin\theta = m(-s\dot{\theta}^2) \tag{C8}$$

From the last two equations, it follows

$$T = m(s\dot{\theta}^{2} + g\sin\theta)$$

$$= \frac{mg}{s} \left[ \frac{2M}{m} D + 2R(1 - \cos\theta) + 3s\sin\theta \right]$$

$$= \frac{2mgR}{s} \left[ \frac{MD}{mR} + (1 - \cos\theta) + \frac{3}{2} \left( \frac{\ell}{R} - \theta \right) \sin\theta \right]$$
(C9)

where Eq. (C5) has been used to obtain the last equality.

We now introduce the function

$$f(\theta) = 1 - \cos \theta + \frac{3}{2} \left( \frac{\ell}{R} - \theta \right) \sin \theta \tag{C10}$$

From the fact  $\ell = (L - D) >> R$ , we may write

$$f(\theta) \approx 1 + \frac{3}{2} \frac{\ell}{R} \sin \theta - \cos \theta = 1 + A \sin(\theta - \phi)$$
 (C11)

where we have introduced

$$A = \sqrt{1 + (\frac{3}{2} \frac{\ell}{R})^2}$$
 ,  $\phi = \tan^{-1} \left(\frac{2R}{3\ell}\right)$  (C12)

From Eq. (C11), the minimum value of  $f(\theta)$  is seen to be given by

$$f_{\min} = 1 - A = 1 - \sqrt{1 + \left(\frac{3}{2} \frac{\ell}{R}\right)^2}$$
 (C13)

Since the tension T remains nonnegative as the particle swings around the rod, we have from Eq. (C9) the inequality

$$\frac{MD}{mR} + f_{\min} = \frac{M(L - \ell)}{mR} + 1 - \sqrt{1 + \left(\frac{3\ell}{2R}\right)^2} \ge 0$$
 (C14)

or

$$\left(\frac{ML}{mR}\right) + 1 \ge \left(\frac{M\ell}{mR}\right) + \sqrt{1 + \left(\frac{3\ell}{2R}\right)^2} \approx \left(\frac{M\ell}{mR}\right) + \left(\frac{3\ell}{2R}\right) \tag{C15}$$

From Eq. (C4), Eq. (C15) may be written as

$$\left(\frac{ML}{mR}\right) + 1 \ge \left(\frac{ML}{mR} + \frac{3L}{2R}\right)(1 - \alpha) \tag{C16}$$

Neglecting terms of the order (R/L) or higher, the last inequality leads to

$$\alpha \ge 1 - \frac{\left(\frac{ML}{mR}\right) + 1}{\left(\frac{ML}{mR} + \frac{3L}{2R}\right)} = \frac{\frac{3L}{2R} - 1}{\frac{ML}{mR} + \frac{3L}{2R}} = \frac{1 - \frac{2R}{3L}}{\frac{2M}{3m} + 1} \approx \frac{1}{1 + \frac{2M}{3m}}$$
 (C17)

The critical value for the ratio D/L is therefore

$$\alpha_c = \frac{1}{1 + \frac{2M}{3m}} \tag{C18}$$

## **Marking Scheme**

# Theoretical Question 1 A Swing with a Falling Weight

T-4-1	C1-	A Swing with a Fairing Weight		
Total Scores	Sub Scores	Marking Scheme for Answers to the Problem		
Part A	(a)	Relation between $\dot{\theta}$ and $\dot{s}$ . $(\dot{s} = -R\dot{\theta})$		
4.3 pts.	0.5	0.5 for proportionantly constant (-rx).		
	(b) Velocity of $Q$ relative to $O$ . $(\bar{v}_Q = R\dot{\theta}\hat{t})$			
	0.5	> 0.2 for magnitude $R\dot{\theta}$ . > 0.3 for direction $\hat{t}$ .		
	(c) Particle's velocity at <i>P</i> relative to <i>Q</i> . $(\vec{v}' = -s\dot{\theta}\hat{r} + \dot{s}\hat{t})$			
	0.7	$\triangleright$ 0.2+0.1 for magnitude and direction of $\hat{r}$ -component.		
	0.7	$\triangleright$ 0.3+0.1 for magnitude and direction of $\hat{t}$ -component.		
(d)		Particle's velocity at <i>P</i> relative to <i>O</i> . $(\vec{v} = \vec{v}' + \vec{v}_Q = -s \dot{\theta} \hat{r})$		
0.7		$\triangleright$ 0.3 for vector addition of $\vec{v}'$ and $\vec{v}_Q$ .		
		$\triangleright$ 0.2+0.2 for magnitude and direction of $\vec{v}$ .		
	(e)	$\hat{t}$ -component of particle's acceleration at $P$ .		
0.7		> 0.3 for relating $\vec{a}$ or $\vec{a} \cdot \hat{t}$ to the velocity in a way that implies $ \vec{a} \cdot \hat{t}  = v^2 / s$ .		
	$ ightharpoonup 0.4  ext{ for } \vec{a} \cdot \hat{t} = -s \dot{\theta}^2  ext{ (0.1 for minus sign.)}$			
	(f)	(f) Potential energy $U$ . > 0.2 for formula $U = -mgh$ .		
0.5 > 0.3 for h		> 0.3 for $h = R(1 - \cos \theta) + s \sin \theta$ or <i>U</i> as a function of $\theta$ , <i>s</i> , and <i>R</i> .		
	(g)	Speed at lowest point $v_m$ . $\triangleright$ 0.2 for lowest point at $\theta = \pi/2$ or $U$ equals minimum $U_m$ .		
	0.7	> 0.2 for total mechanical energy $E = mv_m^2/2 + U_m = 0$ .		
		$\triangleright$ 0.3 for $v_m = \sqrt{-2U_m/m} = \sqrt{2g[R + (L - \pi R/2)]}$ .		
Part B	(h)	Particle's speed $v_s$ when $\overline{QP}$ is shortest.		
4.3 pts.	2.4	$\triangleright$ 0.4 for tension T becomes zero when $\overline{QP}$ is shortest.		
1		> 0.3 for equation of motion $-T + mg \sin \theta = m(-s\dot{\theta}^2)$ .		
		> 0.3 for $E = 0 = m(s\dot{\theta})^2 / 2 - mg[R(1 - \cos\theta) + s\sin\theta]$ .		
		$ ightharpoonup 0.4  ext{ for } \frac{3}{2}(\theta_s - \frac{L}{R}) = \tan\frac{\theta_s}{2}.$		
		$\triangleright$ 0.5 for $\theta_s = 9\pi/8$ .		
		$ > 0.3+0.2 $ for $v_s = \sqrt{4gR/3}\cos\pi/16 = 1.133\sqrt{gR} $		

	(*)			
	(i)	The speed $v_H$ of the particle at its highest point.		
		▶ 0.4 for particle undergoes projectile motion when $\theta \ge \theta_s$ .		
1.9		$\triangleright$ 0.3 for angle of projection $\phi = (3\pi/2 - \theta_s)$ .		
		> 0.3 for $v_H$ is the y-component of its velocity at $\theta = \theta_s$ .		
		<ul> <li>0.4 for noting particle does not strike the surface of the rod.</li> <li>0.3+0.2 for</li> </ul>		
		$v_H = \sqrt{4gR/3}\cos(\pi/16)\sin(\pi/8) = 0.4334\sqrt{gR}$ .		
Part C	(j)	The critical value $\alpha_c$ of the ratio $D/L$ .		
3.4 pts	3.4	> 0.4 for particle's energy $E' = MgD$ when the weight has stopped. > 0.3 for $s = L - D - R\theta$ .		
		> 0.3 for $E' = MgD = mv^2 / 2 - mg[R(1 - \cos \theta) + s \sin \theta]$ .		
		$\rightarrow 0.3 \text{ for } -T + mg \sin \theta = m(-s\dot{\theta}^2).$		
		> 0.3 for concluding T must not be negative.		
		$\triangleright$ 0.6 for an inequality leading to the determination of the range of $D/L$ .		
		> 0.6 for solving the inequality to give the range of $\alpha = D/L$ .		
		$ ightharpoonup 0.6  ext{ for } \alpha_c = (1 + 2M/3m).$		